

Aberration Theory

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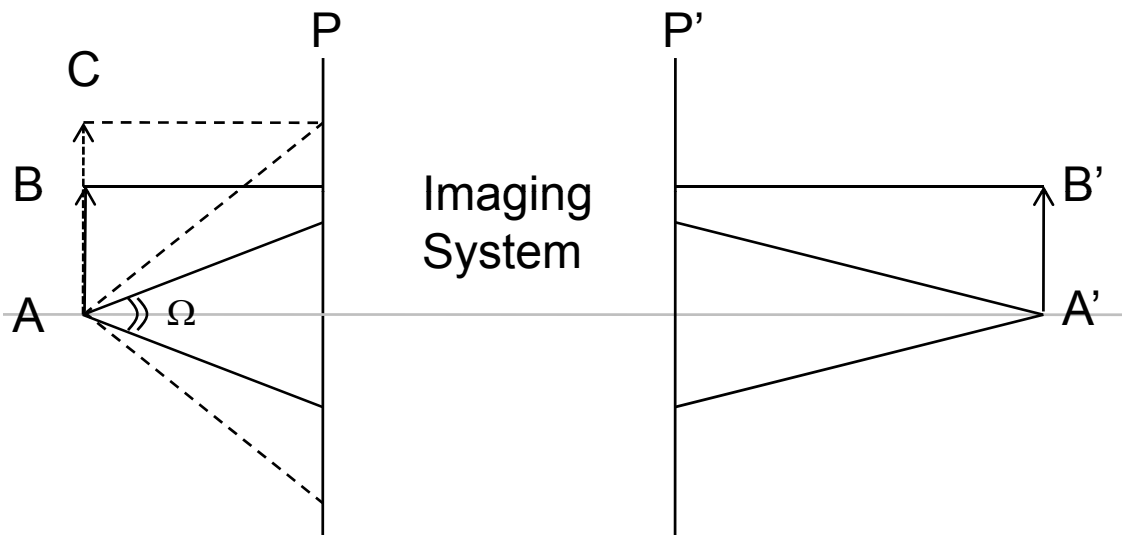
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<http://light.ece.uiuc.edu>



1. Pupils and Windows

- Consider an object imaged by the system below:



! The dotted lines do not make it through the system, i.e, they are blocked somewhere



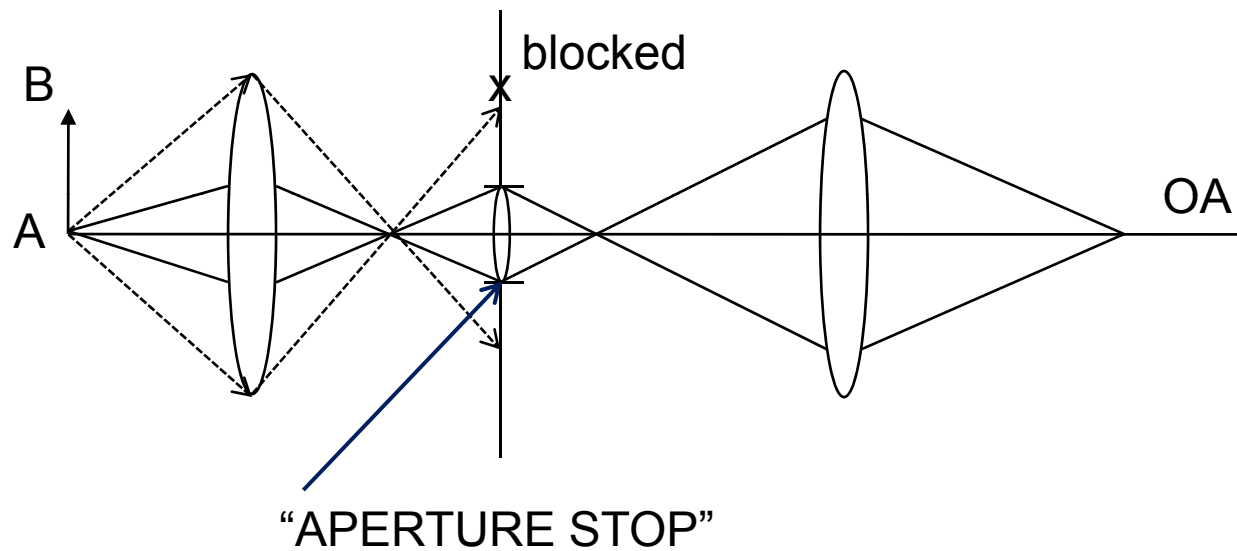
1. Pupils and Windows

- Two main type of limitations:
 - a) Solid angle Ω from the point on axis is limited by the “ENTRANCE PUPIL”
 - b) field of view (i.e extent of the object) limited by “ENTRANCE WINDOW”



1. Pupils and Windows

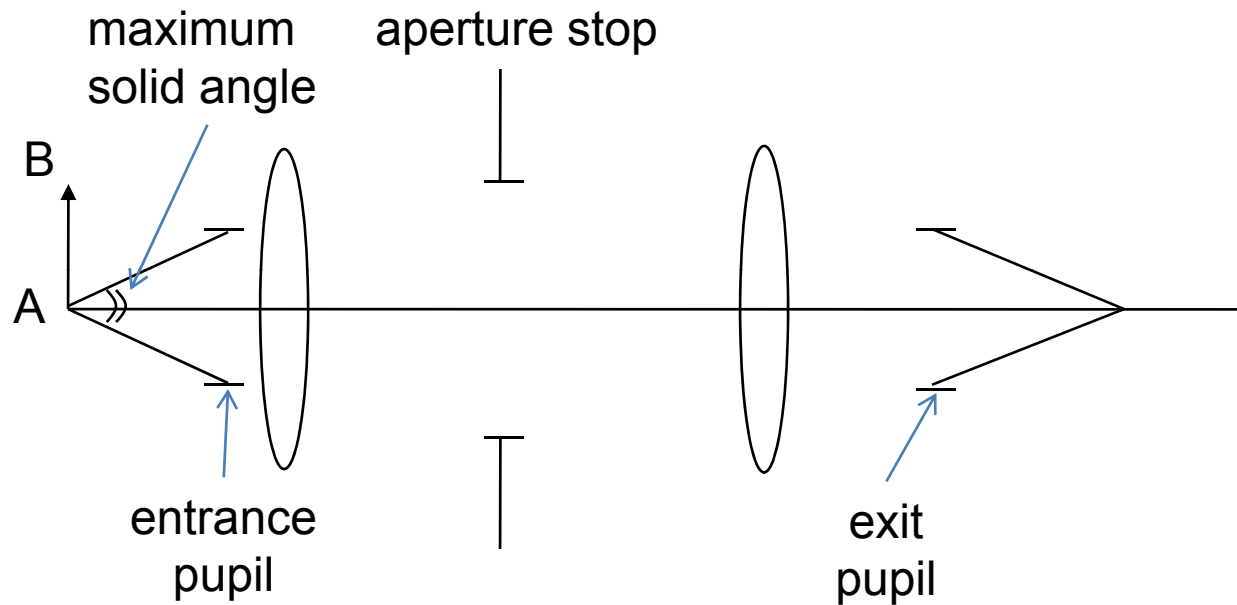
a) Pupils



- Entrance Pupil \equiv image of aperture stop in the object space



1. Pupils and Windows



- Exit Pupil \equiv image of aperture stop in the image space



1. Pupils and Windows

- Entrance Pupil:
 - Limits solid angle Ω
 - Thus, limits the amount of light, i.e brightness of image
 - Angle from object is proportional to spatial frequency

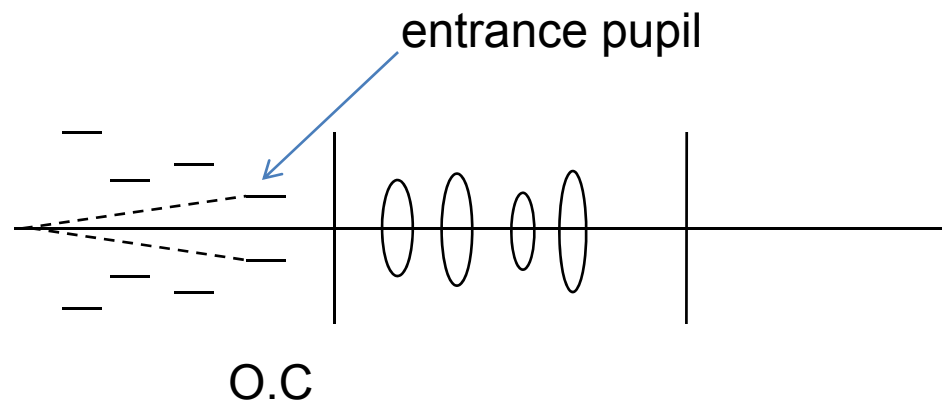
So:

- Adjusting the pupil of our eyes, we adjust brighthness and resolution
- High brightness = high resolution, but may induce saturation and aberations
- Common issue in SLR photography



1. Pupils and Windows

- Note: To find entrance pupil, image all optical elements in object plane and pick the one that subtend the smallest angle from object

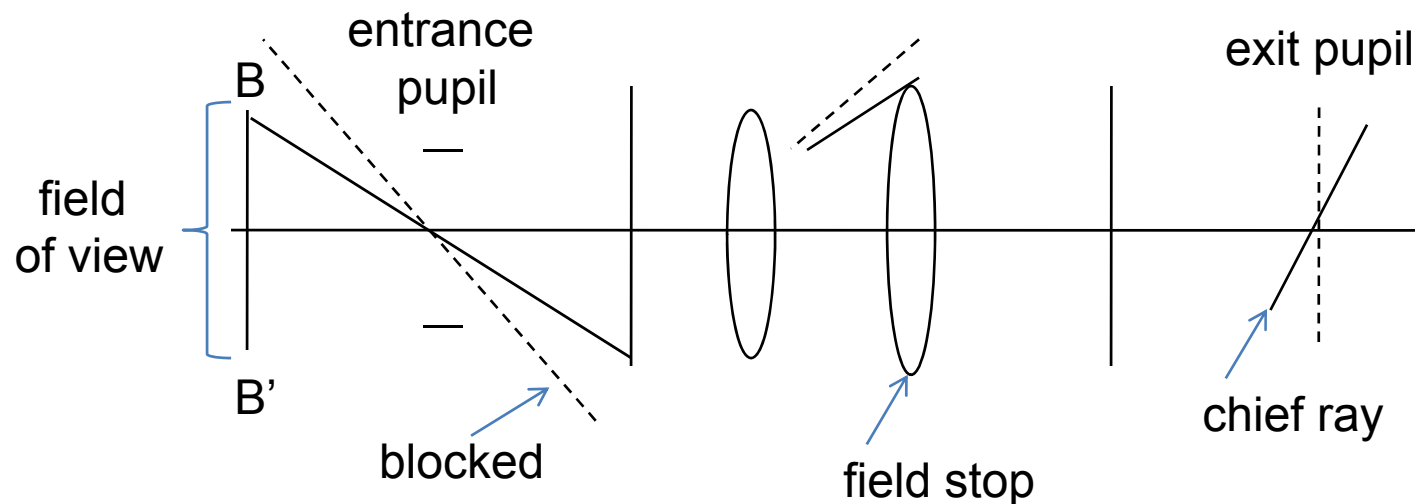




1. Pupils and Windows

b) Windows

- Limit the solid angle made by rays at the center of the entrance pupil (chief rays) with OA

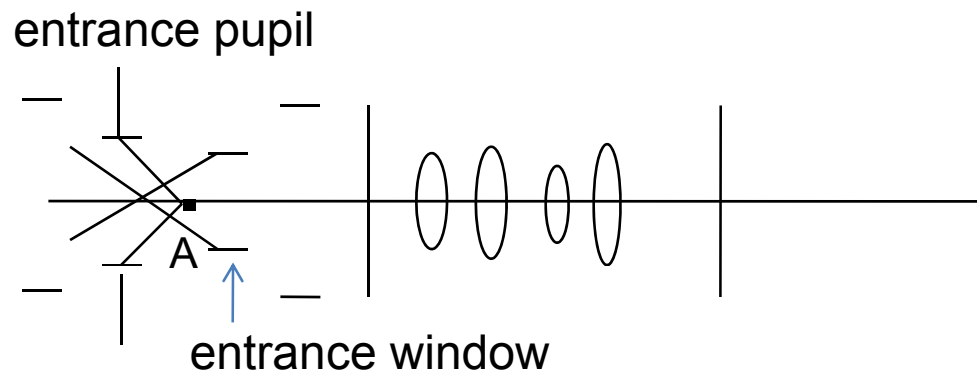


- Entrance Pupil = image of field stop in the object space
- Exit Pupil = image of the field stop in the image space



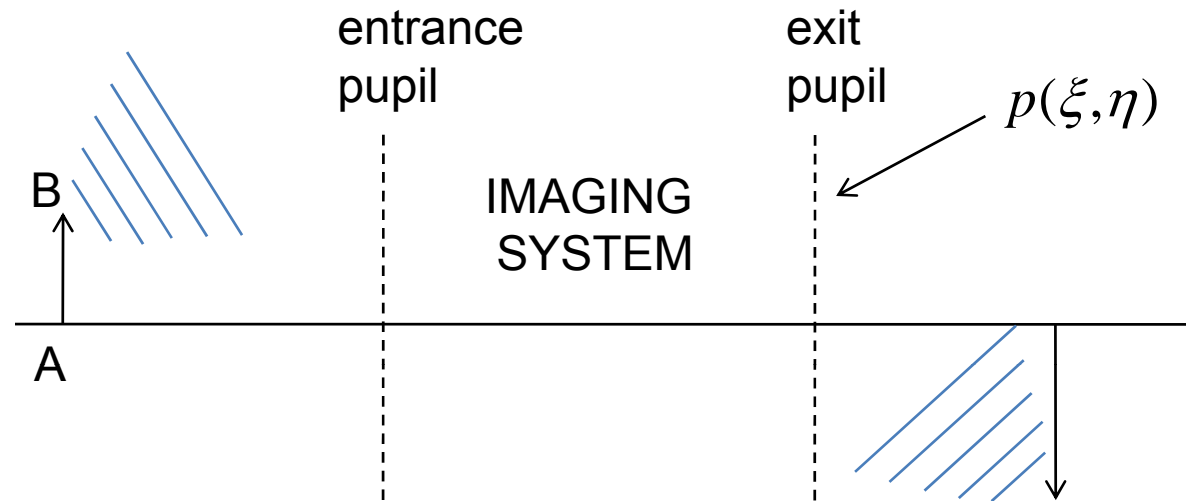
1. Pupils and Windows

- Limits the field of view
- Note:
 - Space and angle are Fourier related
 - It is always a compromise between field of view and solid angle (NA)
 - To find the entrance window: image all components in object space, and pick the one subtending the smallest angle from the center of the entrance pupil.





2. Frequency analysis of coherent imaging



- Recall that the image field equals the object field convolved with the impulse response of the system (Eq 4.4)

$$\bar{U}(x', y') = U(x, y) \otimes h(x, y) \quad 1 \text{ a)}$$

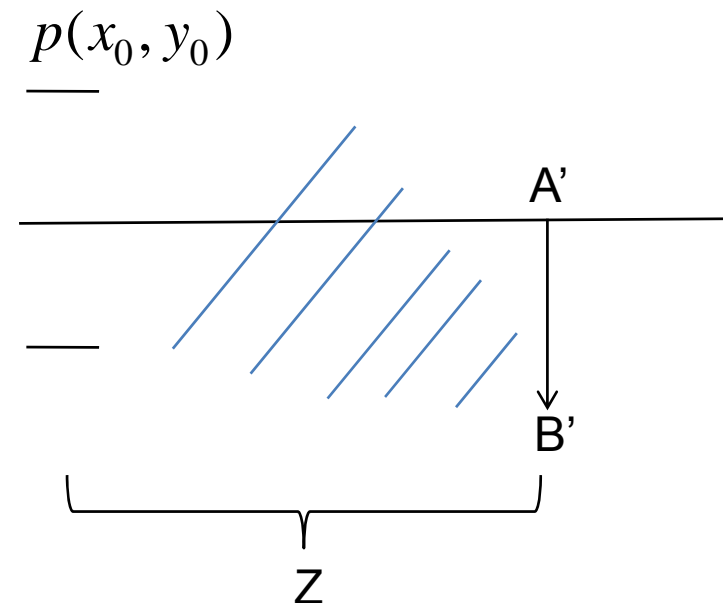
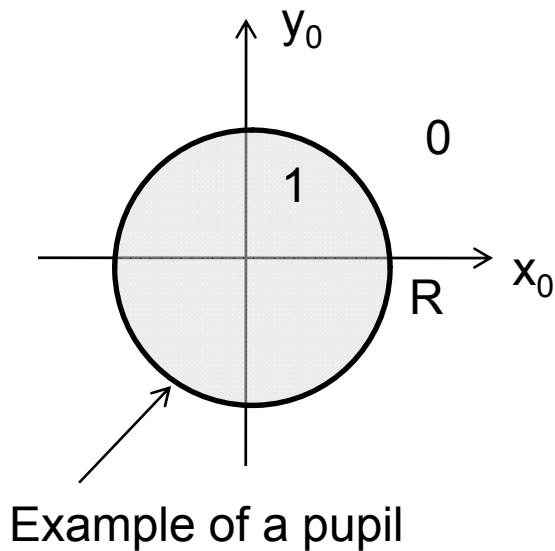
$$h(x, y) = \mathfrak{I}[H(\xi, \eta)] = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} H(\xi, \eta) e^{-i2\pi[x\xi + y\eta]} d\xi d\eta \quad 1 \text{ b)}$$

H = transfer function (coherent)



2. Frequency analysis of coherent imaging

- $p(x_0, y_0)$ = exit pupil
- Note: $\xi = \frac{x_0}{\lambda_z}$; $\eta = \frac{y_0}{\lambda_z} \Rightarrow H(\xi, \eta) = P(\lambda_z \xi, \lambda_z \eta)$





2. Frequency analysis of coherent imaging

- Most common pupil function:

$$p(x_0, y_0) = \begin{cases} 1, & x_0^2 + y_0^2 \leq R^2 \\ 0, & \text{otherwise} \end{cases} \quad (2)$$

- So: Impulse response h is the Fraunhofer diffraction pattern of P !
- This is a “diffraction-limited” instrument, i.e. The best we can do in practice



3. Incoherent Imaging

- If the illuminating field is spatially incoherent, i.e:

$$\langle U(x_1, y_1; t) U^*(x_2, y_2; t) \rangle = I(x_1, y_1) \delta(x_1 - x_2; y_1 - y_2) \quad (3)$$

- The system becomes linear in intensities:

$$\bar{I}(x', y') = I(x, y) \textcircled{v} |h(x, y)|^2 \quad (4) \quad |h|^2 = h \cdot h^*$$

- The intensity impulse response: $|h|^2$!

$$\mathfrak{T}[h \cdot h] \rightarrow \textcircled{v}$$

- Incoherent transfer function:

$$\mathfrak{T}[h \cdot h^*] \rightarrow \otimes$$

$$H_1(\xi, \eta) = \mathfrak{T}[|h(x, y)|^2] = \mathfrak{T}[h \cdot h^*] = H(\xi, \eta) \otimes H^*(\xi, \eta) \quad (5)$$

= autocorrelation of coherent transfer function

- Optical transfer function: $OTF = H_1(\xi, \eta) / H_1(0, 0)$ (5)'

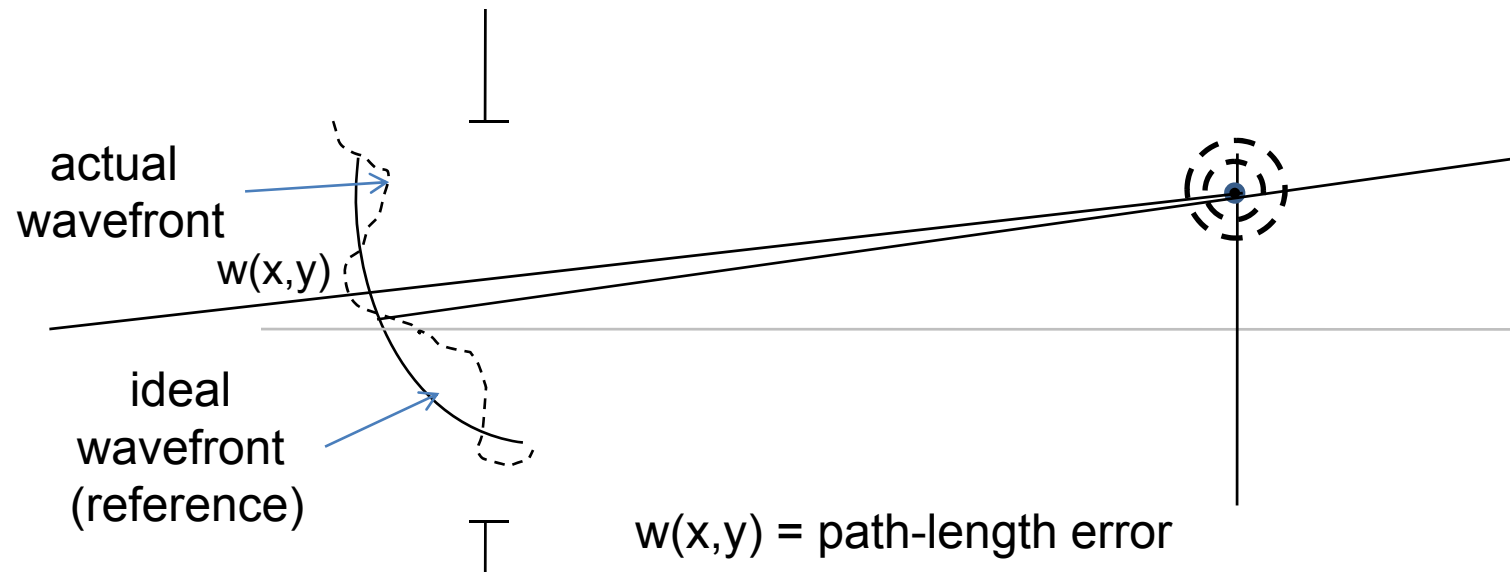


4. Effects of Aberrations

- So far, we assumed that points are imaged into points,
i.e perfect imaging
i.e diffraction-limited system
i.e point-source object generates a spherical wavefront at exit pupil

$$\sin[\theta] \neq \theta$$

- In reality, the outgoing wavefront is nonspherical, i.e aberrated





4. Effects of Aberrations

All the effects of wavefront errors, i.e aberrations can be accounted for by generalizing the pupil functions:

$$P'(x, y) = P(x, y)e^{ikW(x, y)} \quad (6)$$

▪ Note:

- $W(x, y)$ can be positive or negative
- W is the local difference between the actual wavefront and the reference one
- h and H follow from Eq (6):

a) Coherent illumination:

$$\begin{cases} H'(\xi, \eta) = P'(\lambda z \xi, \lambda z \eta) = P e^{ikW(x, y)} \\ h(x, y) = \mathfrak{F}[H(\xi, \eta)] \end{cases} \quad (7)$$



4. Effects of Aberrations

- W can have severe effects on the final image
- Exercise: Calculate h for $k \cdot W = a(\xi^2 + \eta^2)$

b) Incoherent Illumination:

- Recall Eq. 5:

$$\begin{aligned} H_1'(\xi, \eta) &= H(\xi, \eta) \otimes H^*(\xi, \eta) \\ &= P(\lambda z \xi, \lambda z \eta) e^{ikW(\lambda z \xi, \lambda z \eta)} \otimes [\quad]^* \end{aligned} \quad (8)$$

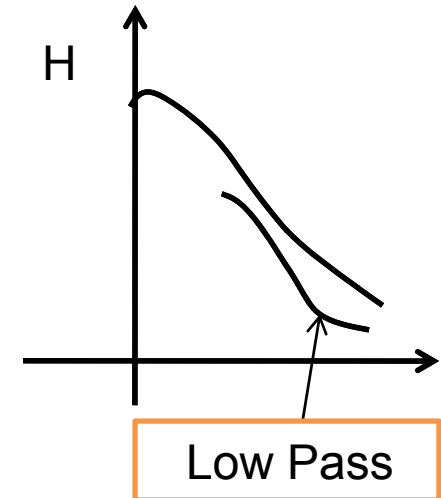
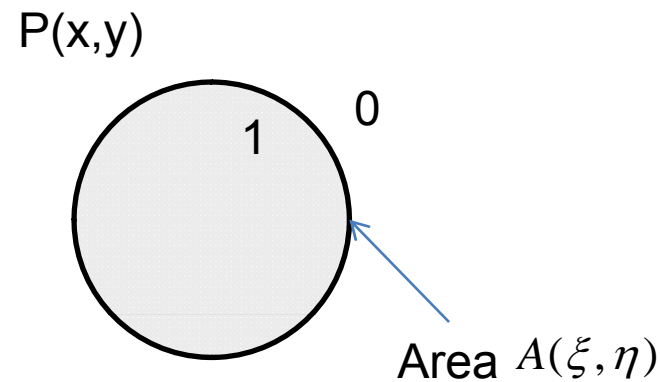
- Expressing the autocorrelation integral for a pupil function

$$\begin{cases} P(x, y) = 1, \text{inside} \\ P(x, y) = 0, \text{outside} \end{cases} \Rightarrow H_1'(\xi, \eta) = \iint_{A(\xi, \eta)} e^{ik \left[W\left(x + \frac{\Delta x}{2}, y + \frac{\Delta y}{2}\right) - W\left(x - \frac{\Delta x}{2}, y - \frac{\Delta y}{2}\right) \right]} dx dy \quad (9)$$



4. Effects of Aberrations

$$\left\{ \begin{array}{l} \xi = \frac{x}{\lambda z} \\ \eta = \frac{y}{\lambda z} \end{array} \right.$$



▪ Note: $|H_1'(\xi, \eta)|^2 \leq |H_1(\xi, \eta)|^2$ (10)

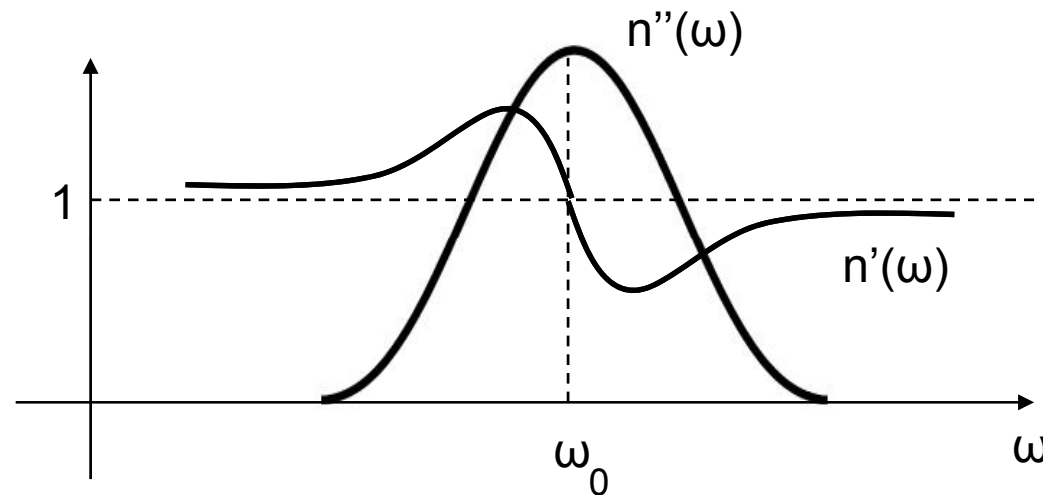
→ Aberrations can only reduce high freq of power spectrum

→ Note that $k=2\pi/\lambda \rightarrow$ wavelength dependence can have an effect, too! Important in broad spectrum imaging (e.g. white light).



5. Chromatic Aberrations

- Recall the dispersion curve from chapter 1.8 (Lorentz model):

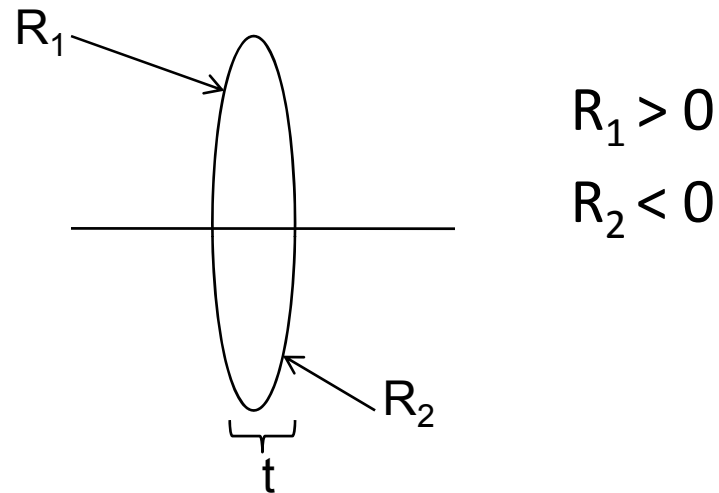


- n' – refractive index
- n'' – absorption
- We assume $n'' \rightarrow 0$, i.e our optical components are highly transparent



5. Chromatic Aberrations

- Consider a lens made of glass with $n'(\omega)$



- Convergence $C = C_1 + C_2 - C_1 C_2 \frac{t}{n}$

$$= \frac{n-1}{R_1} + \frac{1+n}{R_2} - \frac{(n-1)(1-n)t}{nR_1R_2} \quad (11)$$



5. Chromatic Aberrations

- Thin Lens: $t = 0$

$$\Rightarrow C = \frac{1}{f}$$

$$\boxed{\frac{1}{f} = (n - 1) \left(\frac{1}{R_1} - \frac{1}{R_2} \right)} \quad (12)$$

- If $|R_1| = |R_2|$

$$\Rightarrow \boxed{\frac{1}{f} = \frac{2(n - 1)}{R}} \quad (13)$$



5. Chromatic Aberrations

- The dispersion effect is very clear:

$$f(\omega) = \frac{R}{2[n(\omega) - 1]} \Rightarrow \text{focal distance varies with color!} \quad (14)$$

- Differentiate Eq. 14:

$$d \left[\frac{1}{f(\omega)} \right] = d \left[\frac{2(n-1)}{R} \right] \quad (15)$$

$$-\frac{1}{f^2} \frac{df}{d\omega} = \frac{2}{R} \frac{dn}{d\omega} \quad (16)$$

$$\boxed{\frac{df}{d\omega} = -\frac{2f^2}{R} \frac{dn}{d\omega}} \quad (17)$$

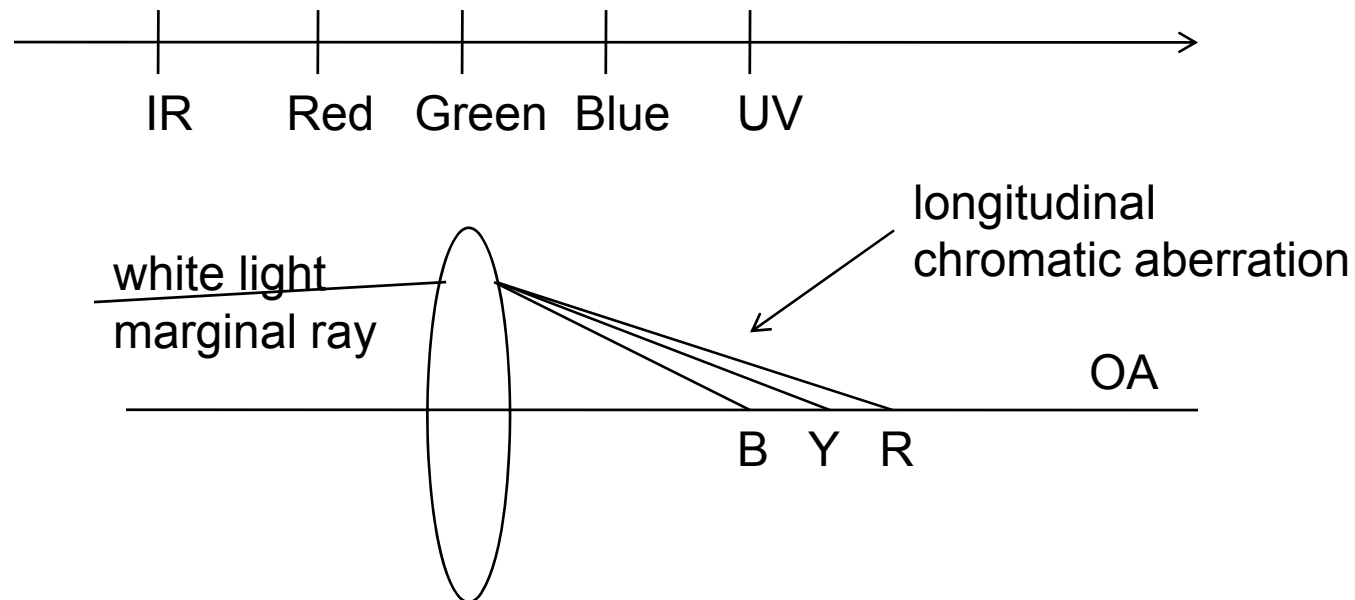


5. Chromatic Aberrations

$$\frac{df}{d\omega} = -\frac{2f^2}{R} \frac{dn}{d\omega}$$

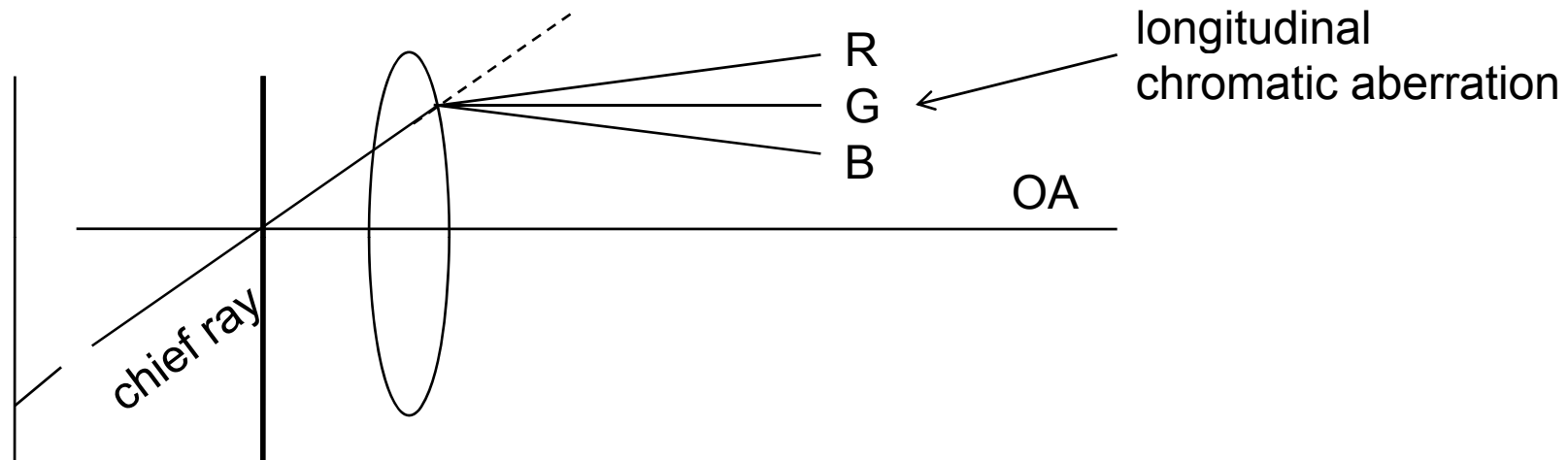
$$\frac{dn}{d\omega} > 0 \Rightarrow \text{normal dispersion}$$

- What is the negative sign telling?





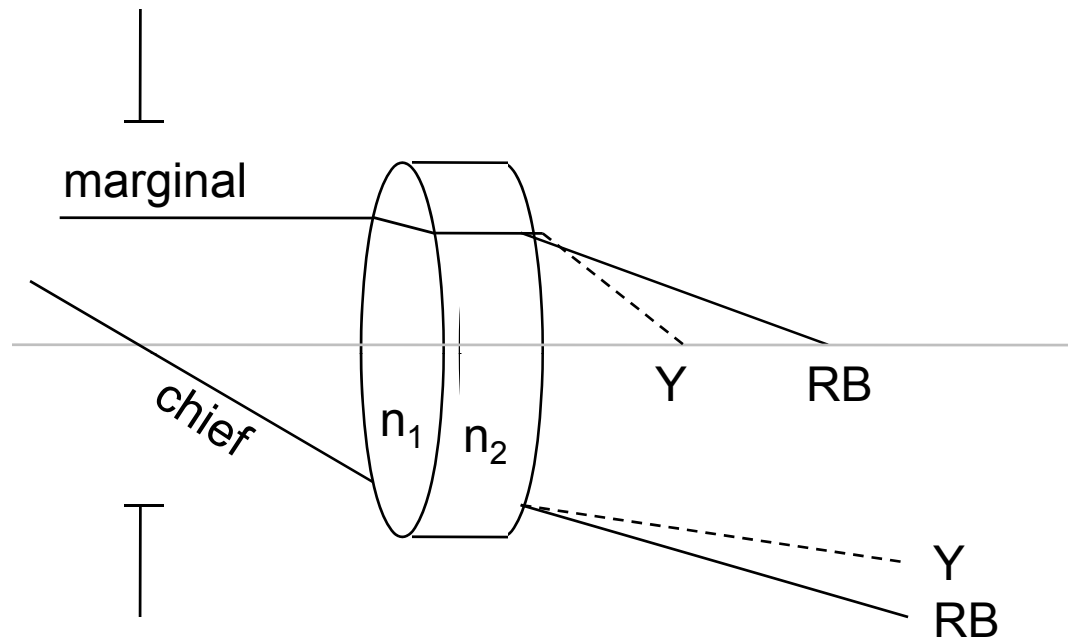
5. Chromatic Aberrations



- Most common correction: the achromatic doublet
 - Sandwich of two lenses, one positive, one negative, which corrects the chromatic aberration at 2 colors (R&B)



5. Chromatic Aberrations



- For thin, closely packed lenses:

$$\frac{1}{f} = \frac{1}{f_1} + \frac{1}{f_2} \quad (18)$$



5. Chromatic Aberrations

- For thin, closely packed lenses:

$$\begin{aligned} \frac{1}{f(\omega_1)} = \frac{1}{f(\omega_2)} &\Rightarrow \frac{2[n_1(\omega_1) - 1]}{R_1} + \frac{2[n_2(\omega_1) - 1]}{R_2} = \\ &= \frac{2[n_1(\omega_2) - 1]}{R_1} + \frac{2[n_2(\omega_2) - 1]}{R_2} \end{aligned} \quad (19)$$

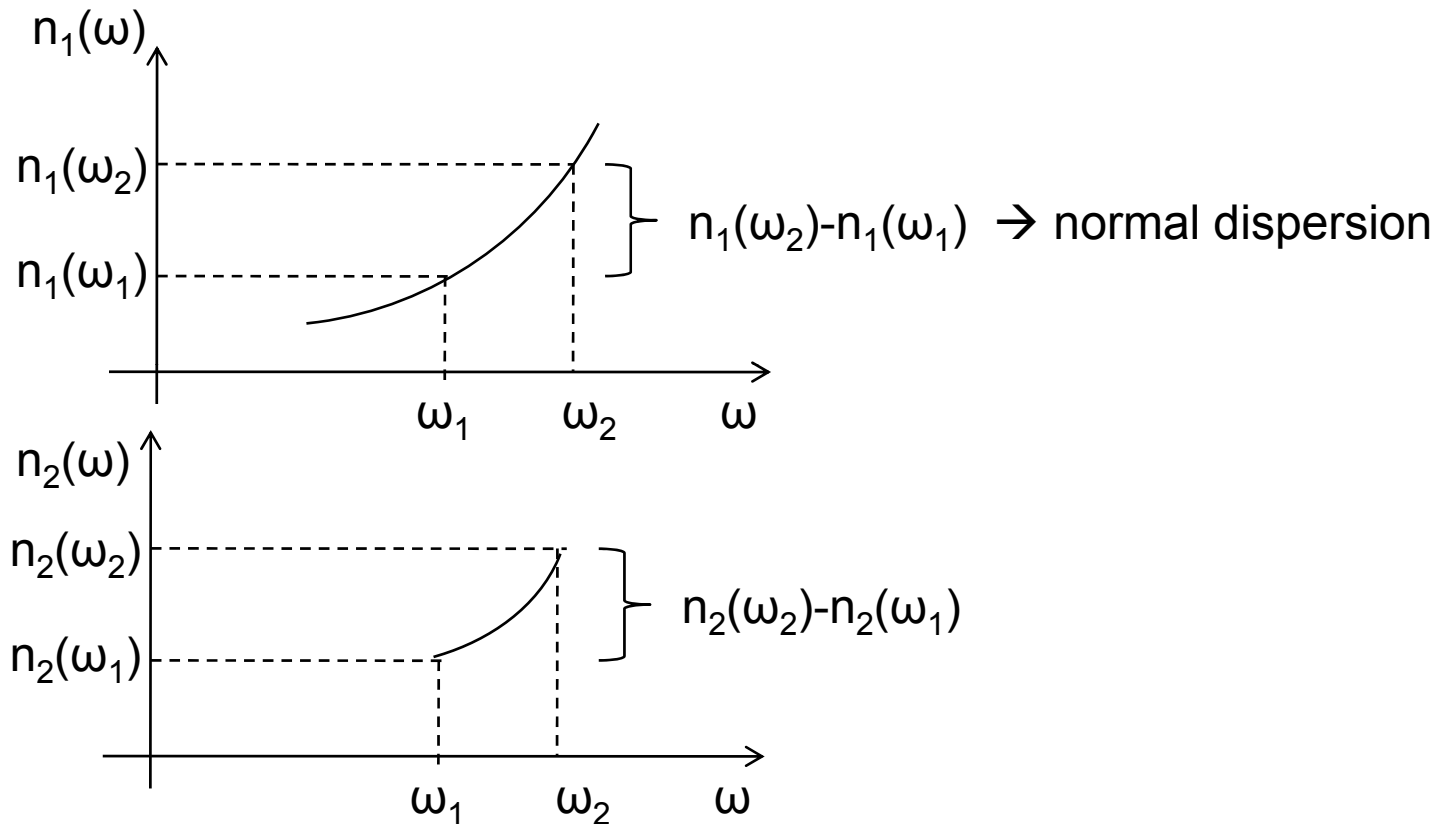
$$\Rightarrow \frac{n_1(\omega_1) - n_1(\omega_2)}{R_1} = -\frac{n_2(\omega_1) - n_2(\omega_2)}{R_2} \quad (20)$$

$R_2 = \text{negative}$

$$\Rightarrow \boxed{\frac{n_2(\omega_1) - n_2(\omega_2)}{n_1(\omega_1) - n_1(\omega_2)} = \frac{R_2}{R_1}}$$



5. Chromatic Aberrations



- This corrects for 2 colors
 - ! A microscope objective may contain more than 8 lenses
 - correction to multiple colors



6. Geometric (monochromatic) aberrations

- $\sin\theta = \theta \rightarrow$ paraxial approximation

$$-\frac{\theta^3}{3!} \rightarrow \text{Third order aberrations (Seidel) = primary aberrations}$$

$$-\frac{\theta^5}{5!} \rightarrow \text{Higher order}$$

- 3rd order aberrations:
 - a) Spherical aberration
 - b) coma
 - c) astigmatism
 - d) field curvature
 - e) distortion

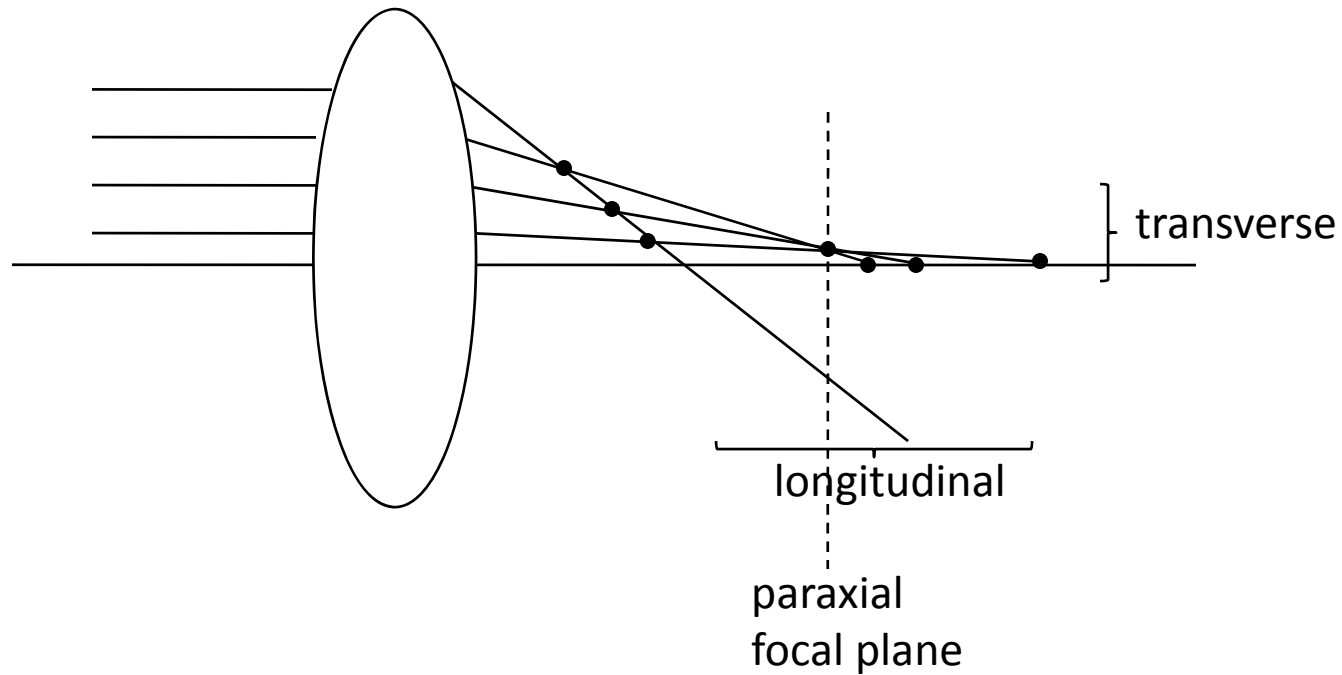
Let's see what each means



6. Geometric (monochromatic) aberrations

a) Spherical aberration

- Focal length depends on aperture





6. Geometric (monochromatic) aberrations

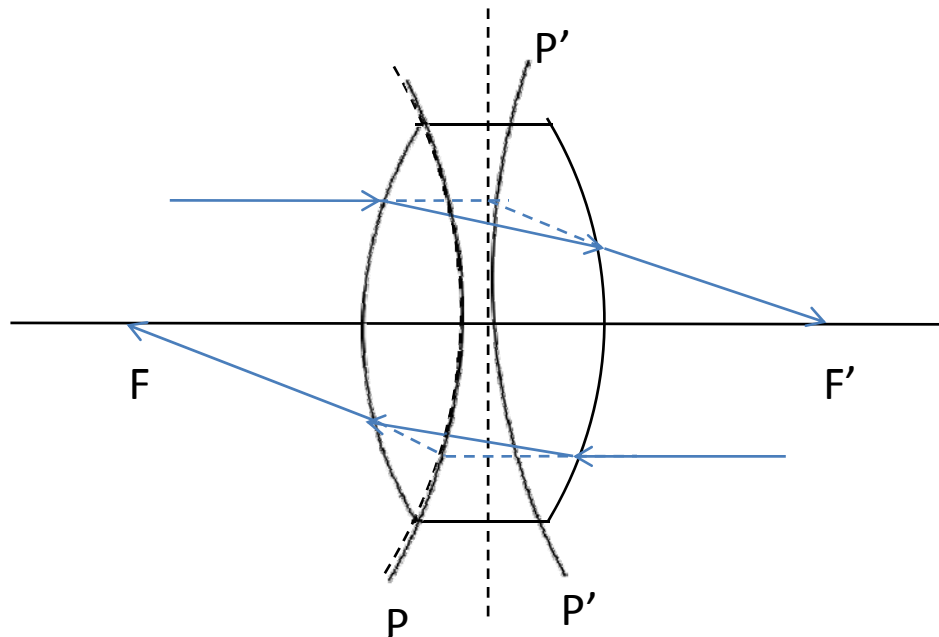
- To avoid SA in practice:
 - Stay close to optical axis
(small aperture)
 - Polish off edge of lens
(as in telescopes, microscope objectives)



6. Geometric (monochromatic) aberrations

b) Coma

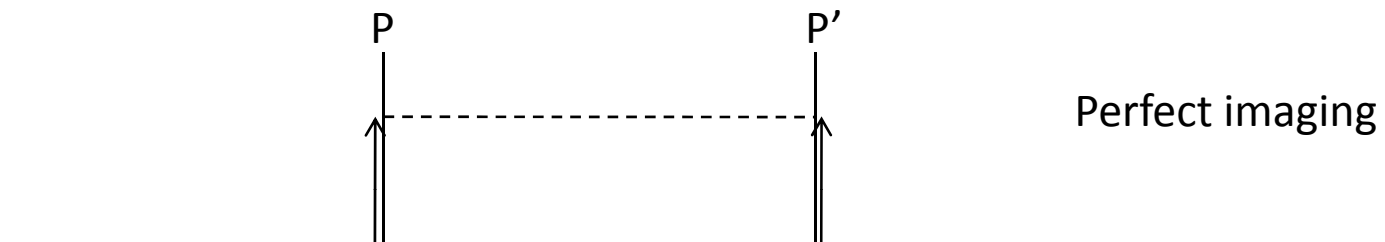
- Occurs because principle planes (i.e. locus of conjugate points image with magnification 1) “curve” at large angles



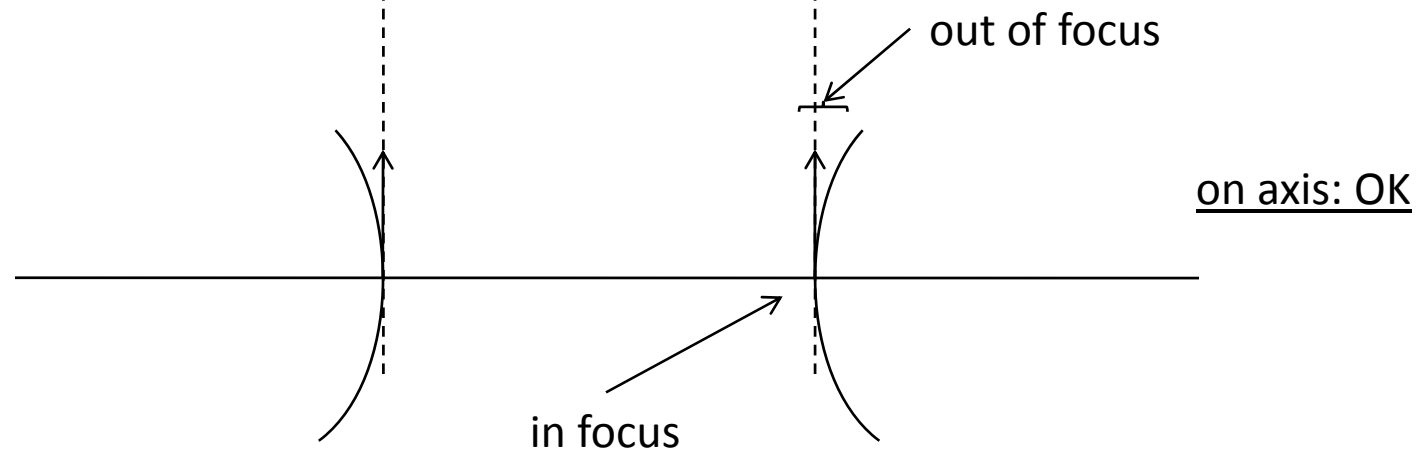


6. Geometric (monochromatic) aberrations

- Paraxial Approximation:

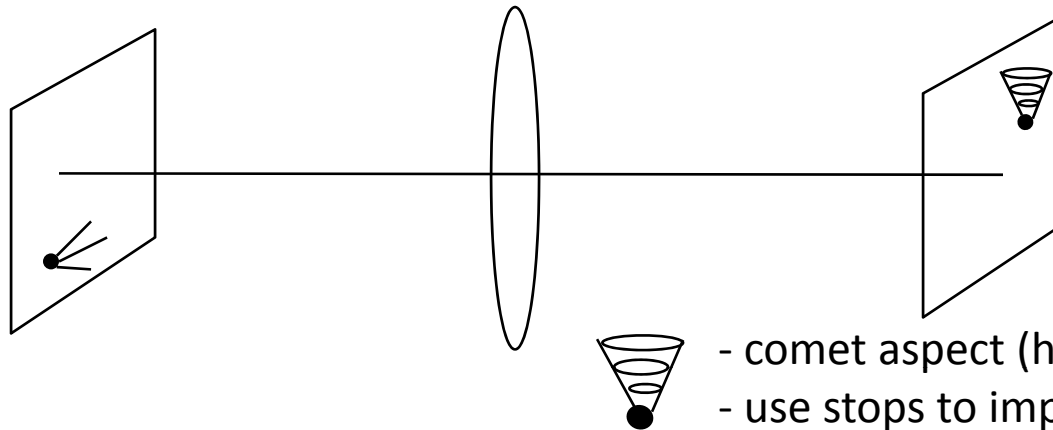


- Real case



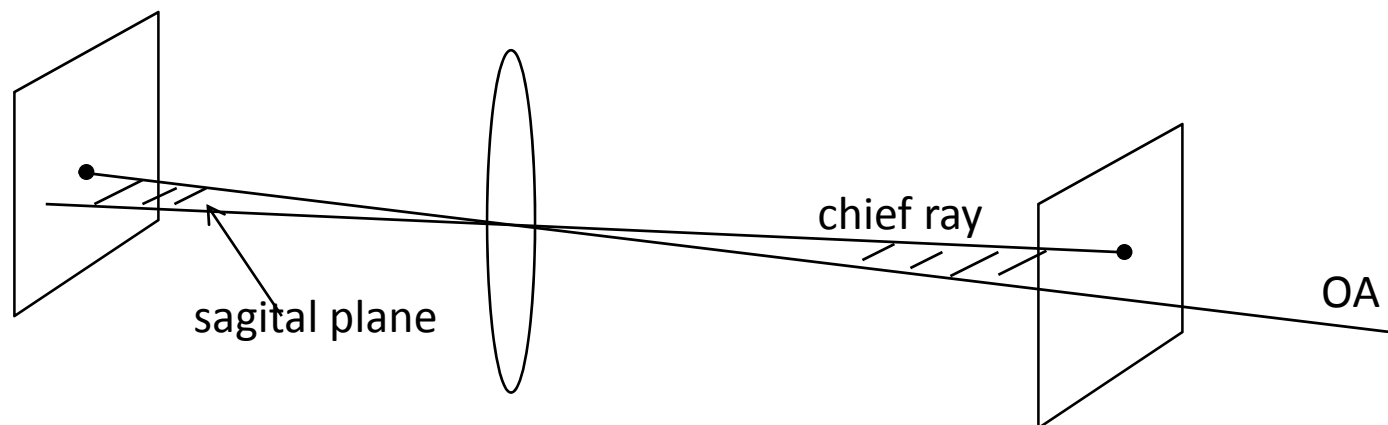


6. Geometric (monochromatic) aberrations



c) Astigmatism

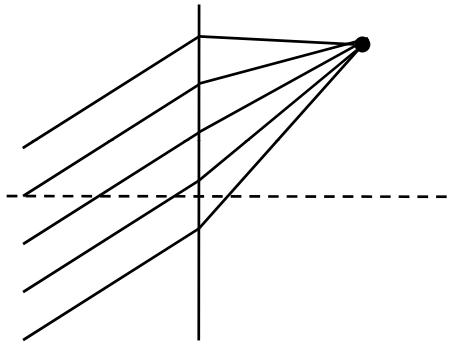
- Different focus on perpendicular planes
- On axis, no difference between the planes



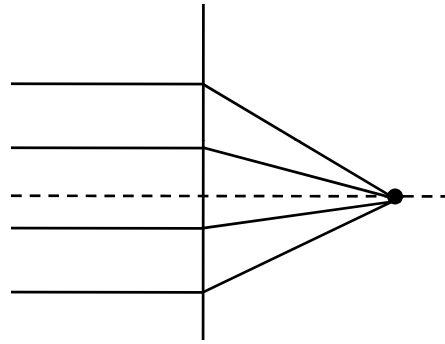


6. Geometric (monochromatic) aberrations

- Meridional plane is perpendicular to sagittal



meridional plane

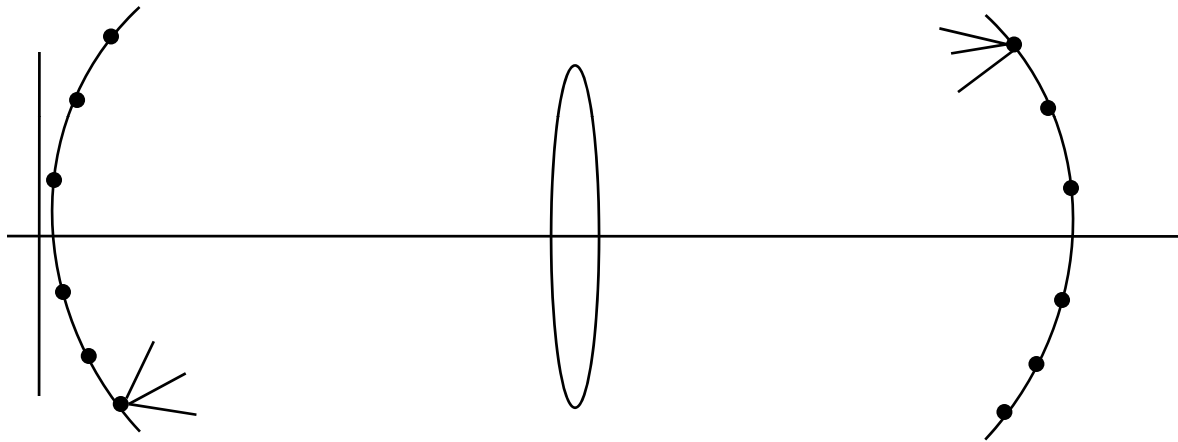


sagittal plane



6. Geometric (monochromatic) aberrations

d) Field curvature



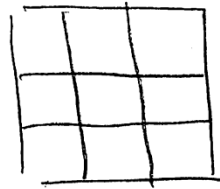
- With astigmatism: 2 paraboloidal surfaces



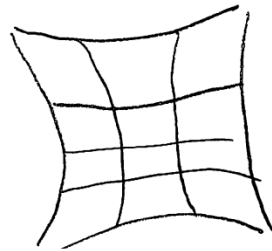
6. Geometric (monochromatic) aberrations

e) Distortion

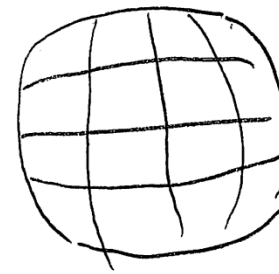
- Transverse magnification is a function of height



a) undistorted



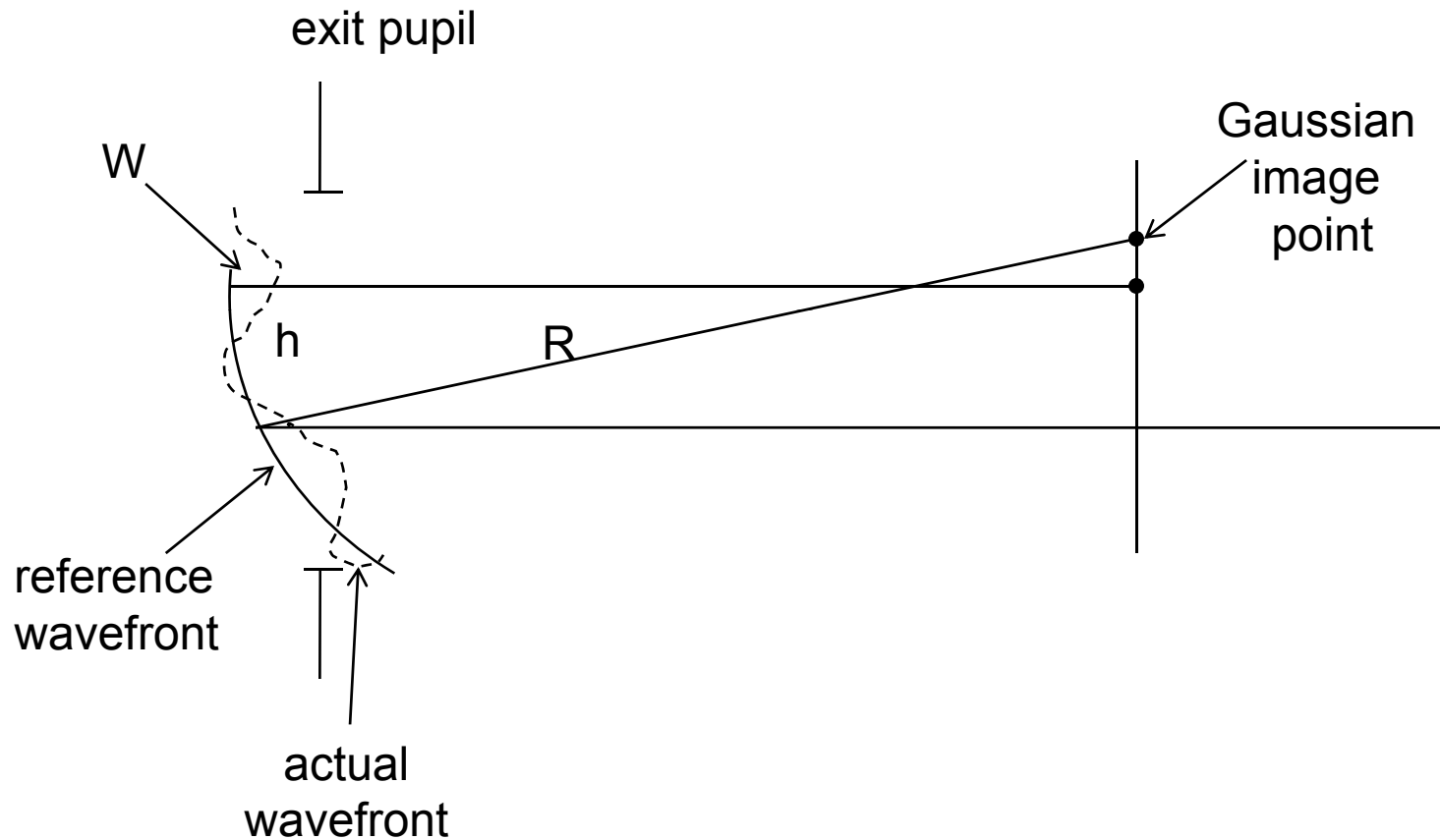
b) cushion
(positive)



c) barrel
(negative)



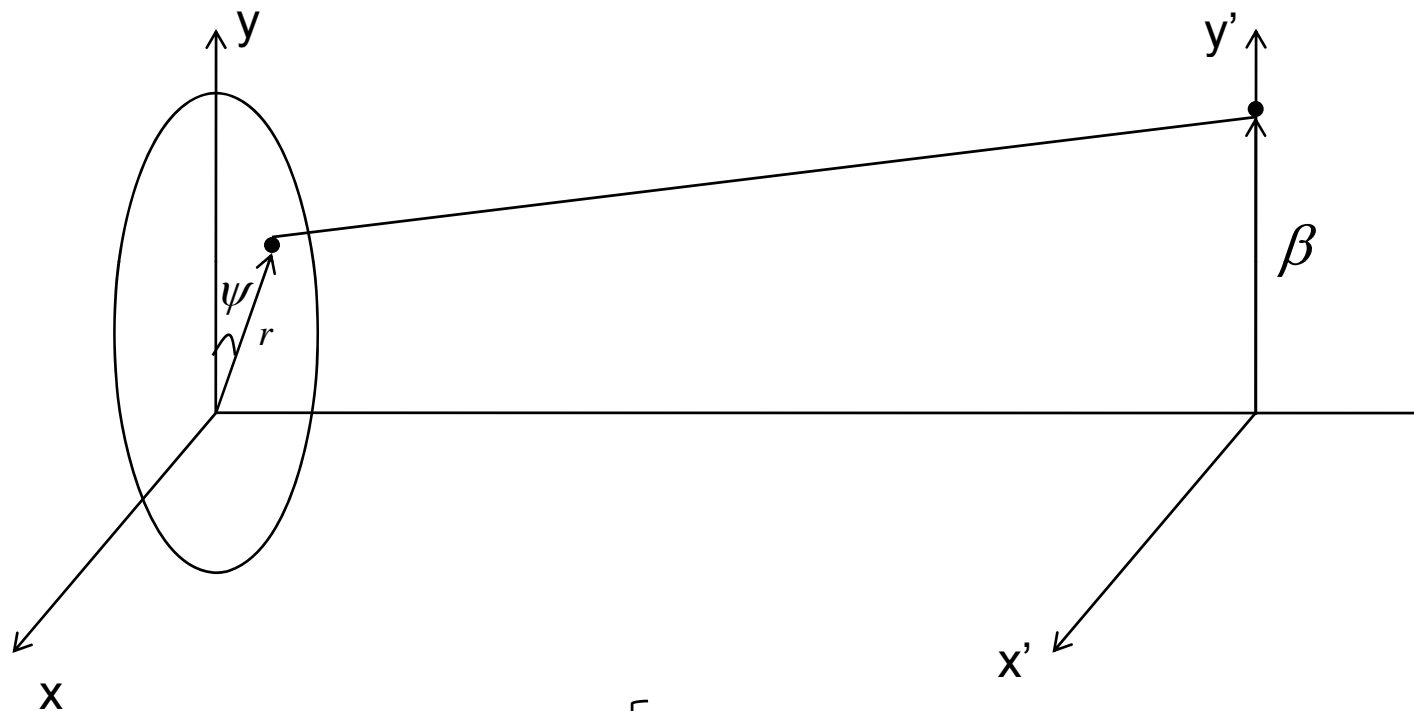
7. Summary of Geometric Aberrations



W = wavefront aberration



7. Summary of Geometric Aberrations



$$W = f(\beta, r, \psi)$$

$$\begin{cases} \beta = \text{normalized field height} \\ r = \text{normalized pupil height} \\ \psi = \text{azimuth angle} \end{cases}$$



7. Summary of Geometric Aberrations

- Symmetries:

$$W(\beta, r, \psi) = W(-\beta, -r, -\psi)$$

→ W depends only on $\beta^2, r^2, \beta r \cos \psi$

If $\beta = 0$ (on axis) → W independent on ψ

→ W depends on $\beta r \cos \psi$

→ $W = W(\beta^2, r^2, \beta r \cos \psi)$



7. Summary of Geometric Aberrations

$$\begin{aligned}
 \Rightarrow W(\beta, r, \psi) = & W_{000} + \\
 & + \cancel{W_{200}}\beta^2 + W_{020}r^2 + W_{111}\beta r \cos \psi + \\
 & \quad \text{defocus} \qquad \qquad \text{tilt} \\
 & + \cancel{W_{400}}\beta^4 + W_{040}r^4 + W_{131}\beta r^3 \cos \psi + \\
 & \quad \text{SA} \qquad \qquad \text{coma} \\
 & + W_{222}\beta^2 r^2 \cos^2 \psi + W_{220}\beta^2 r^2 + \\
 & \quad \text{astigmatism} \qquad \qquad \text{field curvature} \\
 & + W_{211}\beta^2 r \cos \psi + \dots \\
 & \quad \text{distortion}
 \end{aligned}$$